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## COMMENT

# Comment on 'Almost-periodic time observables for bound quantum systems'

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#### Abstract

Hall (2008 *J. Phys. A: Math. Theor.* **41** 255301) makes two claims on a time operator constructed in Galapon (2002 *Proc. R. Soc.* A **458** 2671). We discuss why both claims are wrong.

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In [1], a time operator  $\hat{T}$  is constructed solely from the eigenvectors and eigenvalues of a given Hamiltonian  $\hat{H}$ . The Hamiltonian is assumed to have a discrete non-degenerate spectrum  $E_k$  with the corresponding eigenvectors  $|E_k\rangle$ , k = 1, 2, 3, ...; moreover, the Hamiltonian is semibounded from below such that  $0 < E_1 < E_2 < E_3 < ...$ , and the eigenvalues satisfy  $\sum_{k=1}^{\infty} E_k^{-2} < \infty$ . The  $|E_k\rangle$ 's form a complete set and span the entire Hilbert space  $\mathcal{H}$ , which allows us to write every vector in  $\mathcal{H}$  in the form  $|\psi\rangle = \sum_{k=1}^{\infty} |E_k| |\psi\rangle |E_k\rangle$ . The domain  $\mathcal{D}_H$  of the Hamiltonian  $\hat{H}$  consists of vectors  $|\phi\rangle = \sum_{k=1}^{\infty} a_k |E_k\rangle$  such that  $\sum_{k=1}^{\infty} E_k^2 |a_k|^2 < \infty$ . The time operator  $\hat{T}$  is given by

$$\hat{T} = i\hbar \sum_{j \neq k} \frac{|E_j\rangle \langle E_k|}{E_j - E_k}.$$
(1)

The domain  $\mathcal{D}_T$  of  $\hat{T}$  consists of vectors of the form  $|\varphi\rangle = \sum_{k=1}^{N} b_k |E_k\rangle$  for some finite but otherwise arbitrary positive integer N. The Hamiltonian eigenvectors  $|E_k\rangle$  belong to  $\mathcal{D}_T$ ; this implies that  $\mathcal{D}_T$  is dense or  $\hat{T}$  is densely defined. The operators  $\hat{H}$  and  $\hat{T}$  satisfy the canonical commutation relation  $[\hat{T}, \hat{H}]|\varphi\rangle = i\hbar|\varphi\rangle$  for all  $|\varphi\rangle$  in  $\mathcal{D}_T$  such that  $\sum_k b_k = 0$ ; these vectors comprise the subspace which we referred to as the canonical domain  $\mathcal{D}_c$  of  $\hat{H}$  and  $\hat{T}$ .

Now Hall claims that  $[\hat{T}, \hat{H}]|E_k\rangle = 0$  for every k, and that  $\mathcal{D}_c$  is not dense. These claims are flagrantly erroneous.

To assert that  $[\hat{T}, \hat{H}]|E_k\rangle = 0$  one implies that  $|E_k\rangle$  belongs to the commutator domain of  $\hat{H}$  and  $\hat{T}$ , in particular,  $|E_k\rangle$  belongs to the domain of  $\hat{H}\hat{T}$ . But  $|E_k\rangle$  does not. Since  $|E_k\rangle$ belongs to  $\mathcal{D}_T$  for all  $|E_k\rangle$ 's, we have

$$\hat{T}|E_k\rangle = \sum_{j \neq k} \frac{i\hbar}{(E_j - E_k)} |E_j\rangle,\tag{2}$$

1

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and in order for  $\hat{T}|E_k\rangle$  to belong to the domain of  $\hat{H}$ , it must be that  $\sum_{j \neq k} E_j^2/(E_j - E_k)^2 < \infty$ , which cannot be satisfied because  $E_j^2/(E_j - E_k)^2 \rightarrow 1$  as  $j \rightarrow \infty$  for a fixed k. Hence the expression  $[\hat{T}, \hat{H}]|E_k\rangle$  does not make sense in the Hilbert space. This prevents us from arriving at the contradiction  $\langle E_k | [\hat{T}, \hat{H}] | E_k \rangle = 0 = i\hbar \langle E_k | E_k \rangle$ , which what Hall may be trying to say because we cannot have the equality  $[\hat{T}, \hat{H}]|E_k\rangle = 0$ , even formally.

Also to assert that  $\mathcal{D}_c$  is not dense one implies that there exists a vector  $|\psi\rangle \neq \mathbf{0}$  in the Hilbert space  $\mathcal{H}$  that is orthogonal to all vectors in  $\mathcal{D}_c$ , that is,  $\langle \psi | \varphi \rangle = 0$  for all  $|\varphi\rangle$  in  $\mathcal{D}_c$ ; otherwise,  $\mathcal{D}_c$  is dense. To show that  $\mathcal{D}_c$  is dense, it is sufficient to demonstrate that  $\mathcal{D}_c$ itself has a dense subspace. Let  $\mathcal{D}'_c$  be the linear span of the vectors  $|\psi_{m,n}\rangle = |E_m\rangle - |E_n\rangle$ for all positive integers m, n;  $\mathcal{D}'_c$  is clearly a subspace of  $\mathcal{D}_c$ . Now let  $|\psi\rangle \neq 0$  be in  $\mathcal{H}$  and orthogonal to all  $|\varphi\rangle$  in  $\mathcal{D}'_c$ . Then it must be that  $\langle \psi_{m,n} | \psi \rangle = 0$ , which implies the equality  $\langle E_m | \psi \rangle = \langle E_n | \psi \rangle$  for all m, n. Since  $|\psi\rangle \neq 0$  there is at least an  $n = n_0$ such that  $\langle E_{n_0} | \psi \rangle \neq 0$ . Then  $\langle E_k | \psi \rangle = \langle E_{n_0} | \psi \rangle$  for all k = 1, 2, ...; or  $|\psi\rangle =$  $\sum_{k=1}^{\infty} \langle E_{n_0} | \psi \rangle | E_k \rangle = \langle E_{n_0} | \psi \rangle \sum_{k=1}^{\infty} | E_k \rangle$ , which does not belong to  $\mathcal{H}$ . Then the only way for  $|\psi\rangle$  to be simultaneously in  $\mathcal{H}$  and orthogonal to  $\mathcal{D}_c$  is for  $|\psi\rangle = 0$ . Hence  $\mathcal{D}_c$  is dense. Hall's 'proof' of the non-denseness of  $\mathcal{D}_c$  is a misunderstanding of the definition of a dense subspace.

### References

[1] Galapon E A 2002 Proc. R. Soc. A 487 2671

[2] Hall M J W 2008 J. Phys. A: Math. Theor. 41 255301